

How good is your port ?

By Bohdan Raczynski

Ported enclosures are known to extend the low frequency output of the loudspeaker system by exploiting Helmholtz resonator created by the compliance of the air inside the enclosure and inertance of the air in the port. Acoustic transformer created this way has its own resonant frequency, F_b , at which most (or all, if there was no losses) of the system acoustic output comes from the port. A fairly obvious implication of this is significant velocity associated with the air flow through the port. This in turn, causes all sorts of non-musical noises to be generated by the port and also distortion and acoustic compression. Depending on port geometry, and required low frequency SPL of the system, the issue can be quite significant.

The problem described above belongs to rather complex field of fluid flow. Assuming incompressible flow, some sophisticated FEM programs would be able to model air turbulence and associated vortex shedding in more detail³, but this is well beyond the scope of my article. Fortunately, existing research results, design material and tests results enable us to formulate an approximate macro view of the problem and look at the acoustic impedance of the port under high air velocity. I would like to stress, that the approach I am taking here is a significant simplification of the physics of the problem. However, the resulting model is quite useful and finds confirmation in practical tests¹.

Intuitive approach to port non-linearity

We are all familiar with the need for good carpentry skills when building speaker boxes. Accuracy of joints and sealing the box is essential for proper operation of all types of boxes, be it sealed, vented, passive radiators and so on. Sealed box means exactly this – the air inside the box is trapped and sealed from the external world. There is no parasitic or accidental leakage from the box, so that mathematical model developed for the enclosure continues to be accurate. Consequently, the box Q_b factor is controlled by the designer and NOT by sloppy workmanship.

Vented enclosure also needs to be “air tight”, of course with the exception of the purposefully introduced opening – port. Just as for the sealed box, the cabinet needs to be sealed, so that no accidental air leakage from the box can occur. Properly executed design would include all sorts of seals and gaskets to ensure, that connector boards or drivers themselves do not cause air leakage. Assuming you have your perfect vented box build, you may expect that the SPL curve and input impedance curve will look as on Figure 1. The impedance curve is very familiar and has two characteristic peaks with the dip between them. The dip is located exactly on the box tuning frequency, $F_b = 25\text{Hz}$. My design is a QB3 type with system parameters as shown on Figure 1, and I have also assumed, that $Q_b = Q_p = 1000$, so I have a very low-loss design. Port diameter in this example is 15cm or 6 inches and the input power to the system is 1watt.

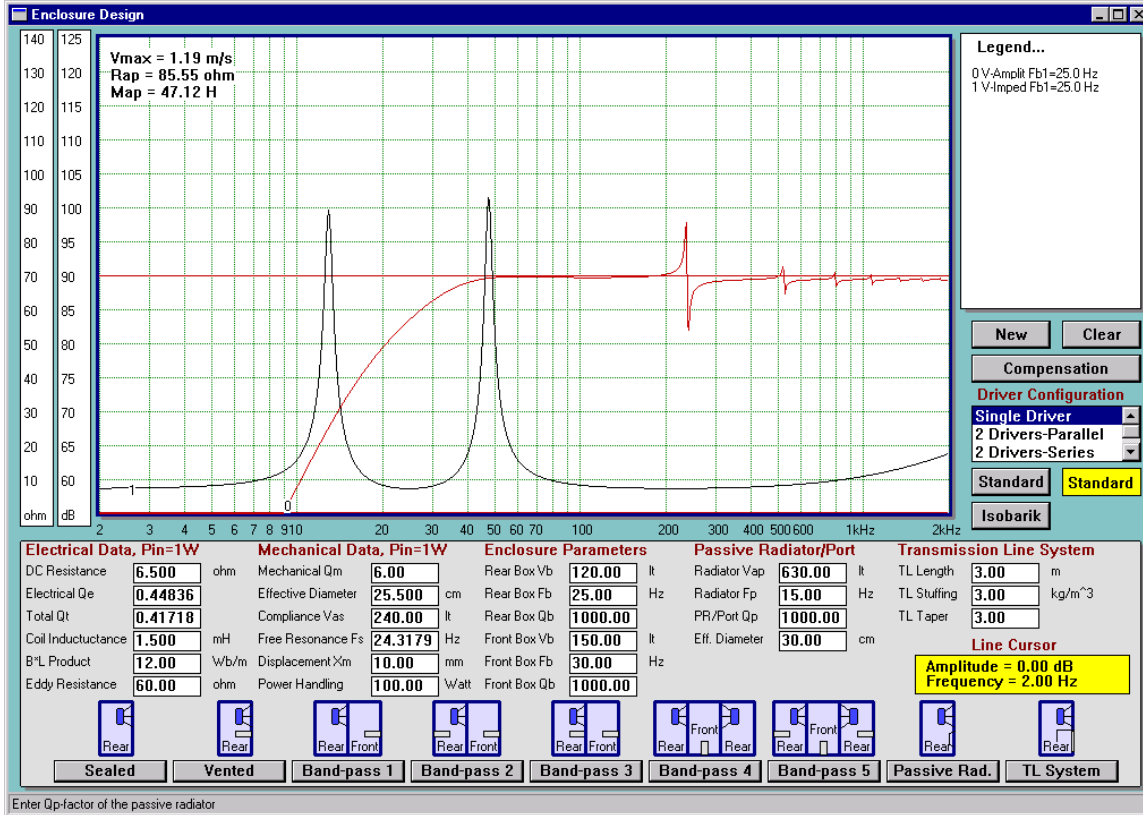


Figure 1 SPL and input impedance @ 1 watt, port diameter = 15cm.

Now, assume, that I start reducing port diameter. Initially, the picture will not change by much. However, when the port eventually becomes very small, intuitively, there should be very little difference between the vented box with a “very small port” and sealed box with a “large leakage” problem⁴. In this case, you would expect that the SPL curve will resemble that of a leaky sealed box and the input impedance curve will lose the lower peak and will become a “single peak” curve just like the sealed boxes have. Between those two extremes, you may expect problems commonly labeled as “port non-linearity”.

Historical research on fluids in tubes

Bies and Wilson⁸ actually experimented with an orifice (more like a decent port) 9cm (3.5”) in diameter and 2.9cm (1.13”) long. They found, that the acoustic resistance, R , of the tube varies with the particle velocity in a similar way as shown on Figure 2

Thurston¹¹, working with fluid flow through circular tubes found that analogous acoustic resistance, R , and inductance, M , again vary in similar way as shown on Figure 2.

A landmark paper by Ingard¹⁰ offers an empirical formula for the non-linear component of the acoustic resistance of a tube. Later on, this work has been expanded by Ingard and Ising⁹ who offered more complete mathematical treatment of orifice behaviour.

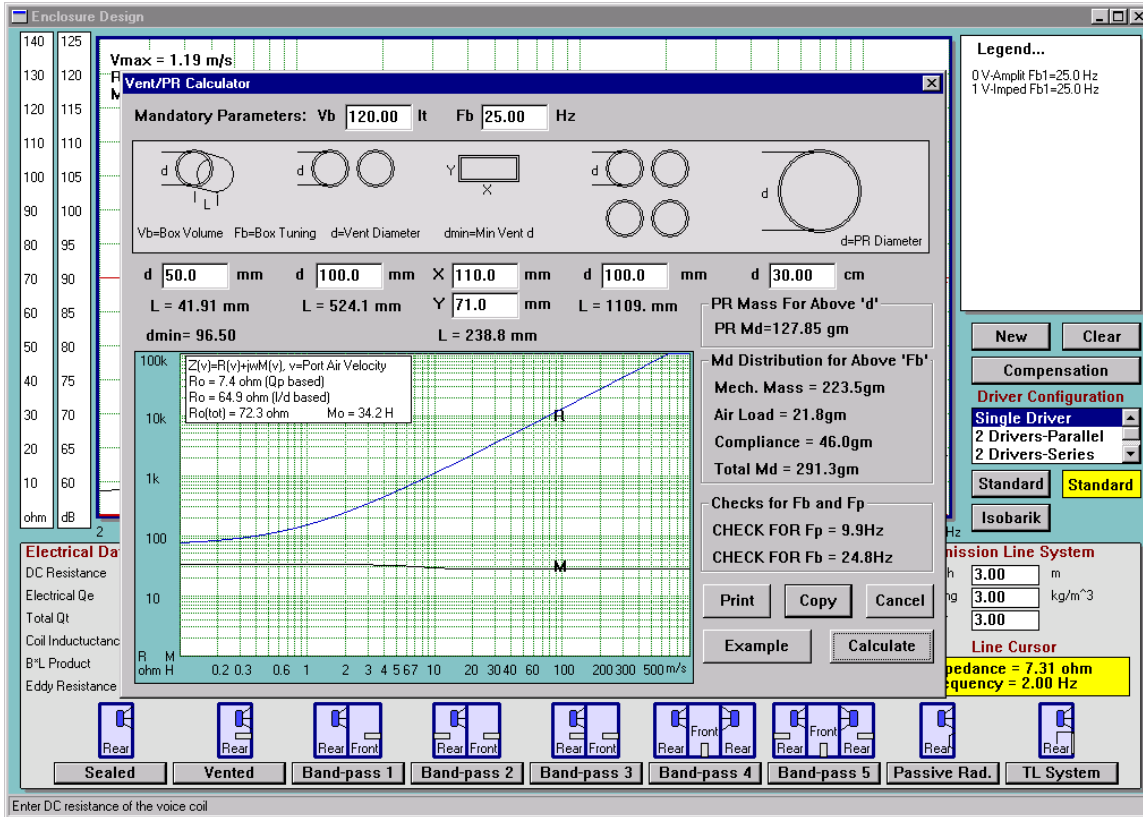


Figure 2. Acoustic impedance $R+j\omega M$ of a short port.

Backman¹, experimenting with port non-linearity, plotted driver's input impedance for various input voltages and determined, that the most sensitive to the amplitude variations is the magnitude of lower impedance peak. He pointed out, that the behaviour of the vented enclosure approaches a sealed enclosure at high levels. The inner diameter of the port was 102mm ($\approx 4''$) and the length of the port was 168mm ($\approx 6.6''$).

Vanderkooy² measured and analyzed ports performance at high levels and proposed a simplified nonlinear model of a port. The comprehensive analysis included jet formation, acoustic compression, turbulence noise and distortion.

Simplified view of the problem

Typically, the area of the cone is many times larger than the area of the port. Since the volume of the air pushed via the port needs to be approximately equal to the volume displacement of the cone, the only way to accomplish this, is to push the air in the port much faster. At low SPL levels, the volume displacement of the cone is small and consequently, the cone velocity is low as well. This, in turn, stipulates low air velocity in the port or what is known – a laminar air flow. The process is characterized by a low Reynolds number, $Re < 2000$, and there is no turbulence in the air flow.

$$Re = (r * V_{max} * \rho) / \mu$$

Where V_{max} is the peak flow rate, ρ is the density of, μ is the viscosity of air, and r is the radius of the port. At this point of time, the acoustic resistance of the port, R_p , (port loss) is linear and is relatively low (50-200ohm).

Also, the mass of air in the port, M_p , is constant. Remember, that mass of the air in the port resonates at F_b (box tuning frequency) with the compliance of the air inside the box, halting the movement of the cone. If there were no losses, at resonance, F_b , all acoustic output of the system would come from the port.

Now, let's start reducing the port diameter. As we can guess, the volume displacement of the cone remains the same, so the air in the port must move faster and faster. Soon, the turbulence occurs in the air flow, jets of air are forming and the process is no longer linear. Reynolds number corresponding to the onset of turbulence is $Re \approx 20000$. When the Reynolds number hits 50000, your vent is compressing. It is now more and more difficult to push the air in the vent. The acoustic output of the port is below the expected level (if the process continued to be linear) and you begin to experience "port compression" effect. Port impedance, Z_p , has now additional component relating to turbulent air flow and depending on the air velocity, $R_p(v)$. Also, the M_p has changed as well, drifting towards 60-70% of its original value and becoming velocity dependant.

The $R_p(v)$ is an interesting element. If plotted in frequency domain, it would resemble the curve depicting port air velocity, shown on Figure 3. This should be of no surprise, as $R_p(v)$ is so heavily dependant on the port air velocity. In the frequency domain, the $R_p(v)$ will quickly reach its peak not far away from the lower peak of the input impedance curve. Depending on the geometry of the port and the volume displacement of the driver, the air velocity in the port may reach 50-100m/s. Such a high values of the air flow result in $R_p(v)$ reaching 5-10kohm levels.

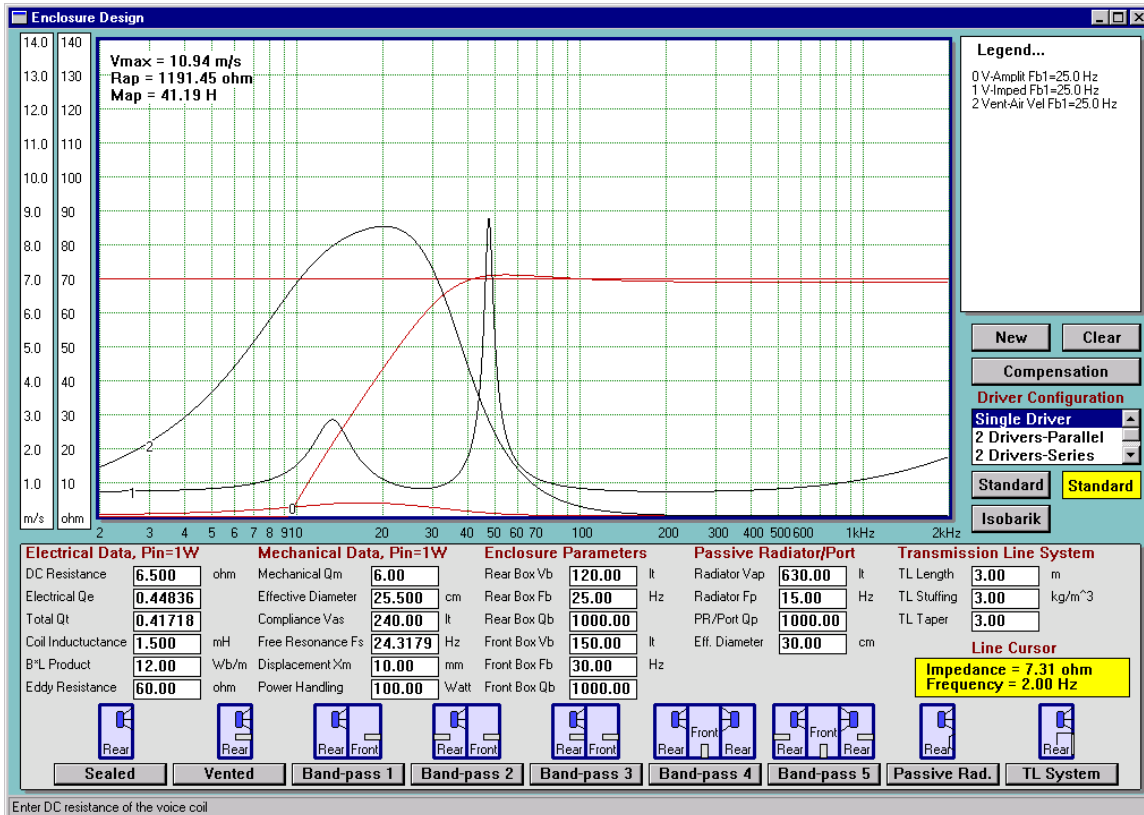


Figure 3. Port air velocity (curve 1) for port diameter = 5cm.

The “double-peak” impedance curve is a clear result of the port action. The port air velocity curve plotted on Figure 3 clearly indicates, that the lower impedance peak will be much more affected by the port non-linearity than the upper impedance peak. Therefore, the $R_p(v)$ needs to be incorporated in the driver’s model.

Changes in driver’s model

Vented enclosure shown of Figure 4, provides different loading for the back of the diaphragm, as compared to the sealed box. The vibrating system and front loading of the diaphragm are represented on the mechanical mobility circuit the same way as for the sealed enclosure.

Introduction of the vent adds several more components such as: (1) mass of the air in the port M_{mp} and its losses R_{mp} and (2) radiation impedance of the port represented by R_{mrp} and M_{mrp} . The air in the port is treated as a mass because of its small volume and more importantly, because it is incompressible. Particles of air will move on both sides of the vent with the same velocity. The air compressed in the box by the back side of the diaphragm has only one path to escape - pushing the air mass through the vent. Therefore, the pressure path consists of series connection of C_{mb} , representing compliant air in the box and the four elements of the port.

Since the air in the port is incompressible, the immediate layer of air in front of the box (radiation impedance) will be connected to the same velocity line as the entry to the port inside the box. The other end of the masses is connected to the $U=0$, or reference velocity as required in mechanical mobility circuits.

The mechanics of the above process can be easily demonstrated on a physical model of a vented box. Connecting a small (1.5V) battery across vented box terminals, we can displace the cone in or out of the box. Small air-flow detecting device (candle) positioned in front of the port will show significant air movements, in the direction opposite to the diaphragm. The volume of air displaced by the cone should be similar to the volume of air leaving the port. If the difference is significant, than leakage losses are significant.

The above experiment clearly shows the pressure (current) path in the mechanical mobility model, so it should now be easy to explain why the compliance of the box is connected in-series with port elements. It is observable, that C_{mb} and $M_{mp}+M_{mrp}$ form series resonant circuit in the mechanical mobility representation. The circuit will act as a “selective short circuit” for the volume velocity U_c , shorting it to $U=0$ (ground) at the circuit resonant frequency. Because of the circuit losses, the short is not perfect, but velocity U_c will be much reduced. In the practical system this situation translates into much reduced cone excursion at the box resonant frequency.

Acoustical impedance representation shows C_{as} and $M_{ap}+M_{arp}$ forming parallel resonant circuit. Electrical circuit theory advocates that very little energy (current) needs to be fed into the circuit for it to resonate and for the current (volume velocity) in the resonant circuit to be still very high.

Therefore, volume velocity in the “feeding” branch, which contains diaphragm output will be very small and volume velocity in the resonant circuit containing port will be high. This effect, although the strongest on the resonant frequency F_b , will extend over some narrow frequency range and on the low-end side creates extended system output. It is the enclosure/port resonance effect, that is being exploited here to augment system output at low frequency. Figure 4 shows mechanical mobility (top diagram) and acoustical impedance (bottom diagram) representation adopted for the vented enclosure model. The components are:

- Cas, equivalent compliance volume Vas transformed to acoustical side.
- Mad, mass of the vibrating system Mms transformed to acoustical side.
- Ras, vibrating assembly loss Rms transformed to acoustical side.
- Mar+Mab, air radiation of the front side of the diaphragm + air load of the back side of the diaphragm.
- Rar, air radiation of the front side of the diaphragm.
- Cab, enclosure compliance Vab transformed to acoustical side.
- Rab, absorption losses of the enclosure transformed to acoustical side.
- Marp, Rarp port radiation.
- Map, mass of the air in the port.
- Rap, frictional losses in the port.

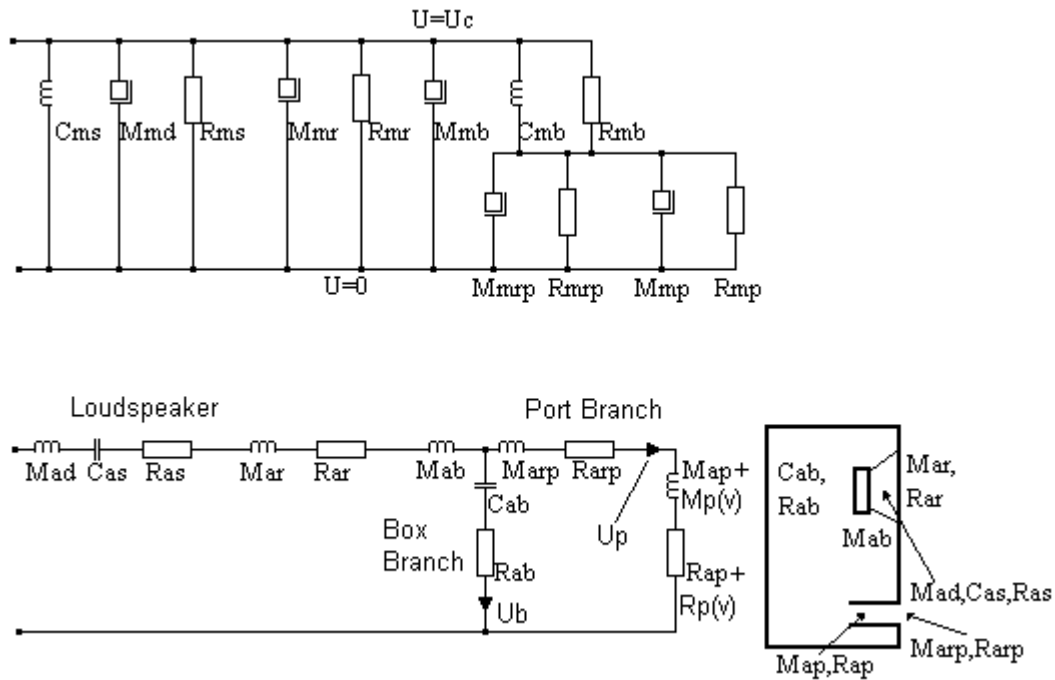


Figure 4. Modified model of the vented box including $Z_p(v) = R_p(v) + j\omega M_p(v)$

The non-linear port impedance was implemented as $Z_p(v) = R_p(v) + j\omega M_p(v)$ in the port branch. Please note, that $M_{ap} + M_p(v)$ will exhibit slight reduction in value as the air speed increases and $R_{ap} + R_p(v)$ will exhibit significant increase in value as the air velocity in the port increases.

Resulting performance

In order to gain some insight into system performance affected by port non-linearity problems, I plotted the SPL for a port of 5cm in diameter for 1W (curve 0), 10W (curve 1) and 100W (curve 2) input power – see Figure 6. As we can see, with the increased input power, there is a sort of “saddle” developing on the SPL curve around the box tuning frequency of 25-30Hz. This is exactly the frequency range, where we would expect the port to contribute most to the system SPL. Our small port is clearly not performing as anticipated.

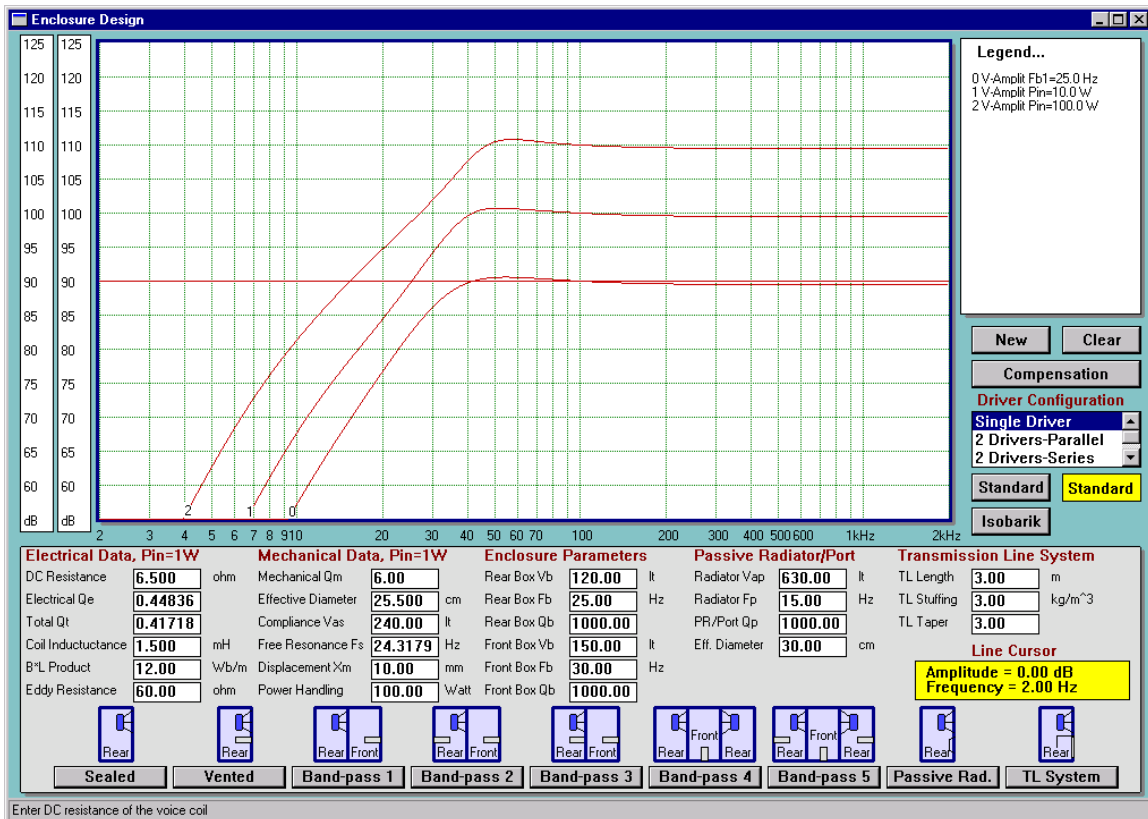


Figure 6. SPL for port diameter of 5cm, @1W, 10W and 100W input power

Next, I modeled SPL for the same power levels, but this time, I used larger port, 15cm in diameter. The resulting plots on Figure 7 do not exhibit the “saddle” any more for 100W power level. Clearly, as the input power is increased, the SPL curves go up, maintaining the approximate shape acquired at 1W power level.

It is also easy to observe, that the SPL curves now have resonant peaks above 200Hz, not seen on Figure 6. This is the result of enlarging the diameter of the port. You may remember, that larger port must also be longer if tuned to the same frequency. Now the length of the port is such, that the self-resonances of the port tube fall into much lower frequency range – just to be displayed on the screen.

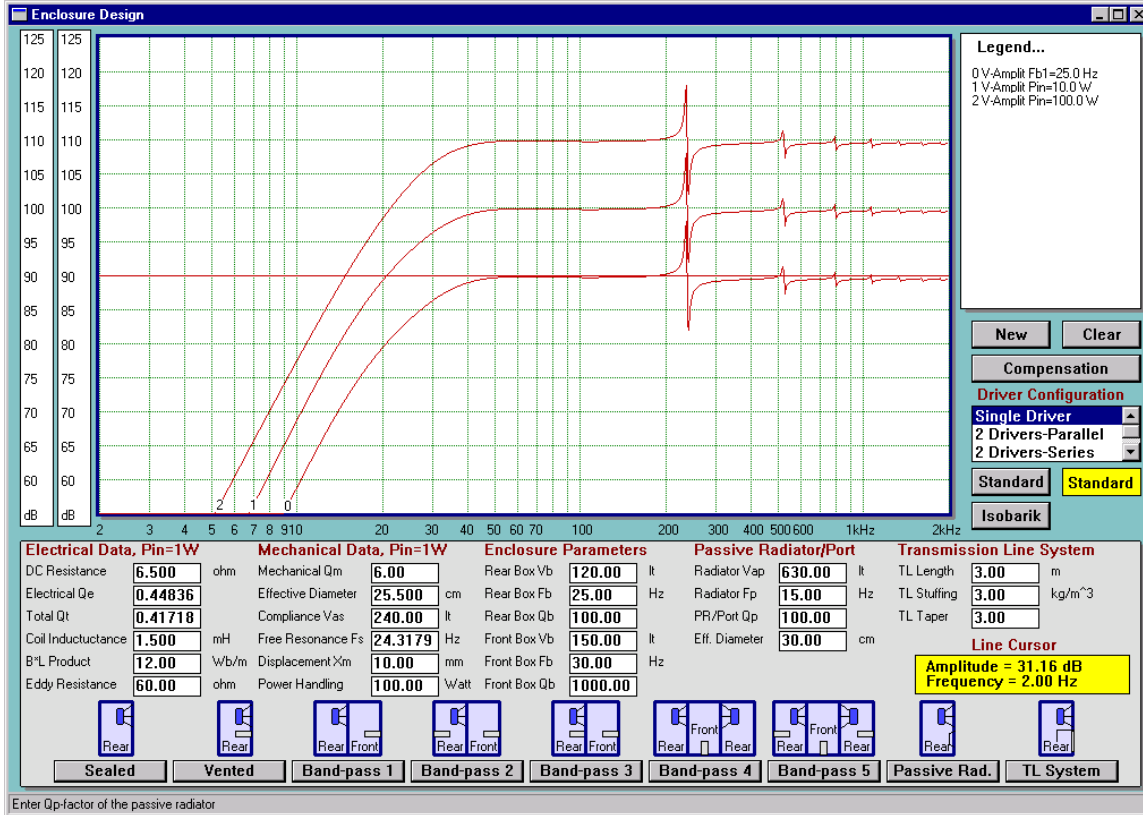


Figure 7. SPL for port diameter of 15cm, @1W, 10W and 100W input power

In order to estimate “port compression” effect, I plotted the SPL of the two ports on the same screen – see Figure 8. Visual inspection of the graphs shows about 5dB of port compression at $F_b = 25\text{Hz}$. Well, I have just about lost my vented box performance gains if I was to use the smaller port.

In the next step, it would be interesting to compare the input impedance plots to see if the lower impedance peak, characteristic for the vented enclosure, indeed disappeared from the plots. The answer is clearly evident on Figure 9, where the input impedance for three ports is plotted at 100W power level. The 15cm vent exhibits the expected “double peaks” curve, however the compression is still registered on the impedance plot. Ideally, the port should still be larger. The Reynolds number calculated for these conditions is $Re = 68000$. Therefore, the port is indeed compressing slightly.

However, the 10cm port has the lower impedance peak significantly reduced. This is a sure sign, that this port is too small for the job.

The severely undersized 5cm port produces “single peak” impedance curve for 100W input power. This port would also prove to be inadequate for 1 watt of input power. The Reynolds number calculated for 1W conditions is $Re = 20000$, which is a clear indication, that the port becomes turbulent. Indeed, the corresponding input impedance plot shown on Figure 10, is a clear indication of the non-linearity problem @ 1 watt for this port.

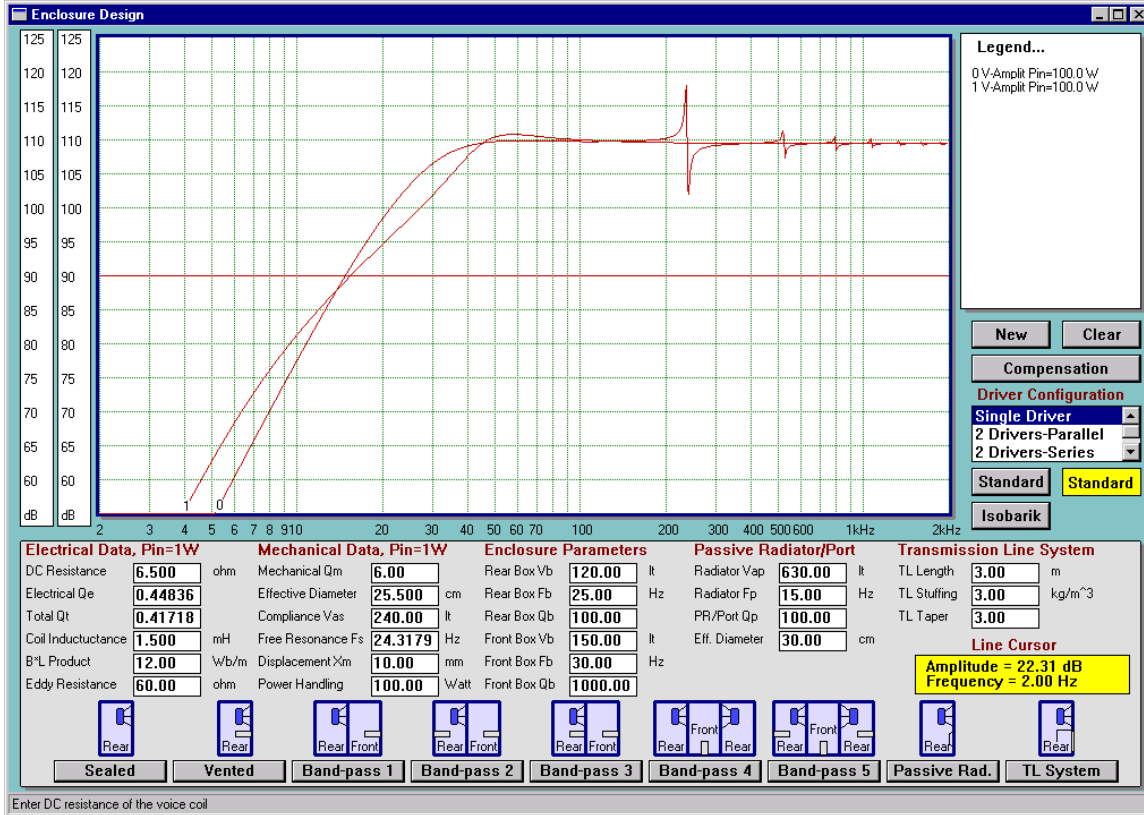


Figure 8. SPL for port diameter of 5cm (curve 0) and 15cm (curve 1) @100W

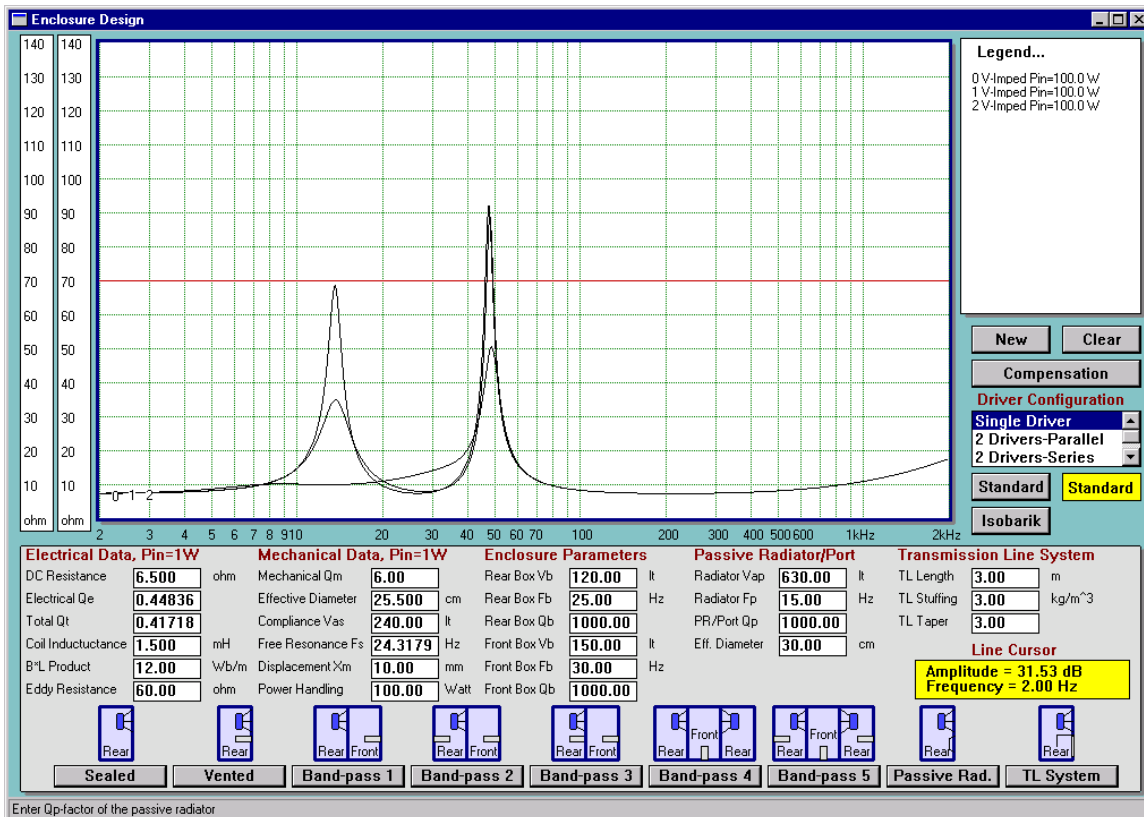


Figure 9. Impedance @100W for port diameter of 15cm,10cm and 5cm.

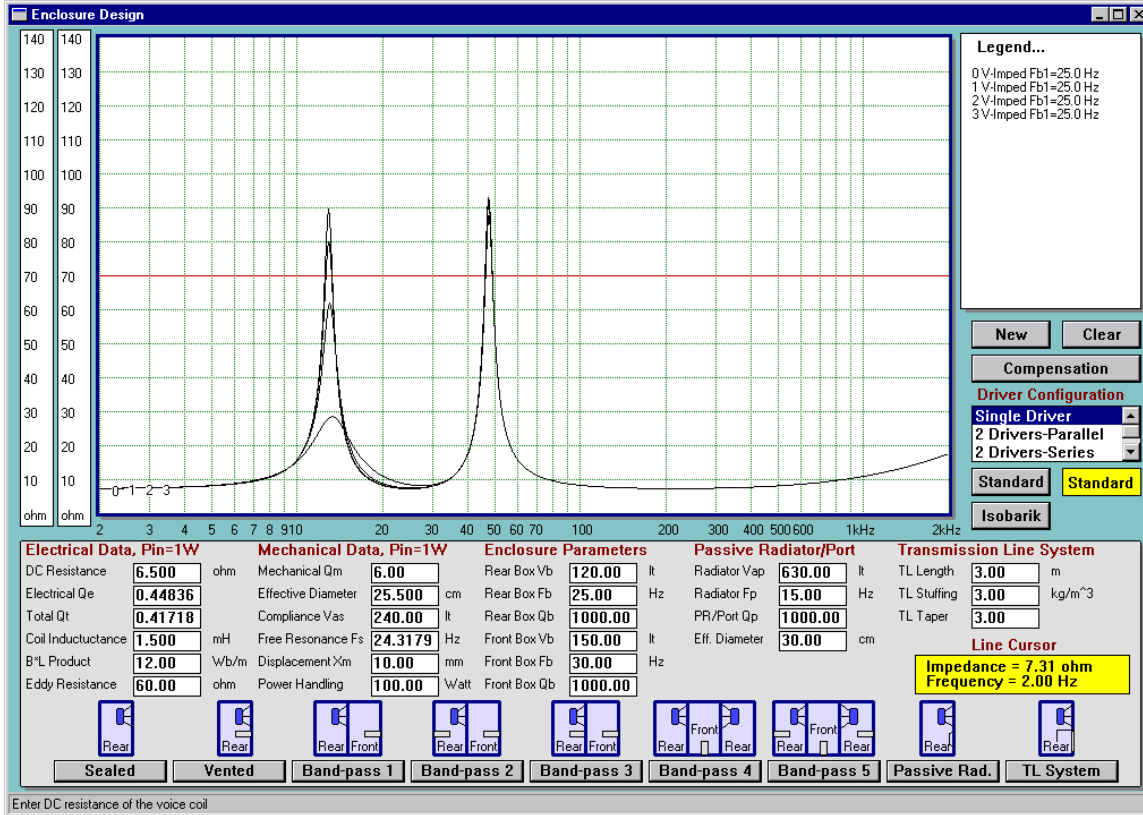


Figure 10. Input impedance @1W for port diameter of 15cm, 10cm, 7.5cm and 5cm.

Remedies

First and foremost, the problem is related to air velocity in the port. As we know¹², air velocity in the port, V_p , depends on volume velocity U_p , via the port branch divided by the port's area, S_p .

$$V_p = U_p/S_p$$

Since the area $S_p = \pi r^2$, it is easy to observe, that air velocity in the port depends on the inverse squared of the port's radius. It seems rather obvious, that port radius should be kept as large as possible.

Dickason¹³ recommends 15cm (6") ports for 15" woofer as a minimum and 10cm (4") ports as a minimum for 12" woofers.

Salvatti⁴ offers a number of recommendations relating to port geometry. This includes balancing inlet and exit flows by using different tapers for inlet and exit.

Roozen³ advocates port geometry based on 6 degrees diverging contour towards the ends of the port. The inlet and outlet are rounded with relatively small curvature.

Conclusions

A simple modification of acoustic impedance of the port in the vented enclosure has been described. The modification consists of adding extra components such as $M_p(v)$ and $R_p(v)$ which are dependant on air velocity inside the port. The modification reflects acoustic compression of the port and shows the changes to the input impedance of the driver under high SPL levels. The changes in input impedance reflect experimental findings¹.

Looking at the performance of the 15cm (6") vent, I was surprised to see the compression evident at 100Watt input power. The vent seems quite large and to tune it to 25Hz in 120liter box, I would have to make it nearly 60cm (23.6") long. This could be a construction challenge, and remember, this vent is still not completely linear at high SPL.

This brief analysis of the vent performance clearly indicates, that nearly all practical size vents will compress at higher SPL levels. Loudspeakers intended for high power stage applications should be give very careful consideration regarding port non-linearity. For this type of application, port compression problem will tend to further degrade system SPL, already affected by the thermal compression during prolonged stage performance.

References

1. J. BACKMAN, The Nonlinear Behaviour of Reflex Ports. AES preprint 3999.
2. J. VANDERKOOY, Nonlinearities in Loudspeaker Ports. AES preprint 4748.
3. N.B ROOZEN, J.E.M. VEAL, J.A.M. NIEUWENDIJK, Reduction of Bass-Reflex Port Nonlinearities by Optimizing the Port Geometry, AES preprint 4661.
4. A. SALVATTI, D. BUTTON, A. DEVANTIER, Maximizing Performance from Loudspeaker Ports, AES preprint 4855.
5. G. TRURSTON and others, Nonlinear Properties of Circular Orifices, JASA, Volume 29, Number 9, 1957.
6. A. NOLLE, Small-Signal Impedance of Short Tubes, JASA, Volume 25, Number 1, 1953.
7. G. THURSTON, C.MARTIN, Periodic Flow through Circular Orifices, JASA, Volume 25, Number 1, 1953.
8. D. BIES, O. WILSON, Acoustic Impedance of a Helmholtz Resonator at Very High Amplitude, JASA, Volume 29, Number 6, 1957
9. U. INGARD, H.ISING, Acoustic Nonlinearity of an Orifice, JASA, Volume 42, Number 1, 1967.
10. U. INGARD, On the Theory and Design of Acoustic Resonators, , JASA, Volume 25, Number 6, 1953.
11. G. THURSTON, Periodic Fluid Flow through Circular Tubes, JASA, Volume 24, Number 6, 1952.
12. R. SMALL, Direct Radiator Loudspeaker Analysis, An Anthology of Articles on Loudspeakers from JAES Vol. 1 -Vol. 25 (1953-1977)
13. V. DICKASON, Loudspeaker Design Cookbook, 5th Edition, Old Colony Sound Lab. 1997.

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